

# LECTURE 1

## 1. INTRODUCTION

In these lecture series we will discuss **ancient solutions** to **parabolic equations** and in particular to **geometric flows**.

By a **geometric flow** we typically mean a non-linear parabolic equation involving the change of a geometric quantity (often the curvature) on a manifold.

The simplest parabolic equation is the **heat equation** which describes the distribution of temperature. In fact, the temperature  $u(x, t)$  at the point  $x \in \mathbb{R}^n$  at time  $t > 0$ , then it satisfies the equation

$$u_t = \Delta u$$

where  $\Delta u = \sum_{i=1}^n u_{x_i x_i}$ .

Two of the remarkable properties of the heat equation are: the **smoothing effect** and **infinite speed of propagation**. Also, another very fundamental property is the **parabolic scaling**: if  $u$  solves the heat equation, so does

$$\hat{u} = \beta u(\alpha x, \alpha^2 t)$$

for any  $\alpha, \beta \in \mathbb{R} \setminus \{0\}$ . All these can be seen from its fundamental solution (also a self-similar solution)

$$\Gamma(x, t) = \frac{1}{(4\pi t)^{n/2}} e^{-|x|^2/4t}, \quad t > 0$$

which has an initial data the **dirac mass**  $\delta_0(x)$  at the origin.

It is easy to see that  $\bar{u}(x, t) := u(x, -t)$  doesn't satisfy the same equation that  $u(x, t)$  but rather the **backward heat equation**

$$\bar{u}_t = -\Delta \bar{u}$$

This is in contrast to other time dependent equations such as the **wave equation**

$$u_{tt} = \Delta u$$

where  $\bar{u}$  satisfies the same equation as  $u$ , i.e. **time can be reversed**.

**Definition 1.1.** A solution  $u(\cdot, t)$  to a parabolic equation is called **ancient** if it is defined for all time  $-\infty < t < T$ ,  $T \leq +\infty$ .

In the special case where the solution  $u(\cdot, t)$  to a parabolic equation is defined for all  $-\infty < t < +\infty$ , it is called **eternal**.

*Remark 1.2.* Since time is cannot not reversed in a parabolic equation one expects that ancient solutions are **very special** and that one may expect to be able to provide a classification of such solutions under certain conditions. Note than in the case of the **wave equation ancient** is equivalent to **globally defined**, hence one doesn't expect that ancient solutions are particularly special !

We will see that **ancient** solutions arise as blow up limits at a **singularities**, hence their classification is important to understand the singularities of a parabolic nonlinear equation, in particular to geometric flows.

**Examples of ancient solutions for the heat equation:**

- i. If  $u(x, t) = U(x)$ ,  $-\infty < t < +\infty$ , where  $U(x), x \in \mathbb{R}^n$  is **harmonic**, then  $u(x, t)$  is in particular ancient (actually eternal).
- ii. Let  $u(x, t) = e^{x_1+t}$ ,  $x = (x_1, \dots, x_n)$ ,  $-\infty < t < +\infty$  is ancient (actually eternal). This is a **traveling wave** solution.
- iii. Let

$$u(x, t) = (4\pi(T - t))^{-n/2} e^{\frac{|x|^2}{4(T-t)}}, \quad -\infty < t < T.$$

Then,  $u(x, t)$  is an ancient solution which is defined up  $T$ .

**1.1. Outline of lectures.** In this series of lectures I will discuss classification results for ancient solutions to parabolic equations and geometric flows. The outline of my lectures is as follows:

- Lecture 1:** Liouville type theorems for ancient solutions to the heat equation on complete non-compact manifolds with nonnegative Ricci curvature.
- Lecture 2:** Classification of ancient solutions to semilinear heat equation related to the blow up analysis of the equation  $u_t = \Delta u + u^p$  on  $\mathbb{R}^n$ , for exponents  $1 < p < \frac{n-2}{n+2}$ .
- Lecture 3:** Ancient compact solutions to curve shortening flow and mean curvature flow.
- Lecture 4:** Ancient compact solutions to the 2-dimensional Ricci flow.

## 2. LIOUVILLE'S THEOREM FOR THE HEAT EQUATION

Let  $M$  be a complete non-compact Riemannian manifold of dimension  $n \geq 2$  with  $\text{Ricci}(M) \geq 0$ . Here and later  $\text{Ricci}(M)$  is the Ricci curvature and a manifold is complete if every geodesic extends to infinity.

**Theorem 2.1** (Yau 1975). *Any **positive harmonic** function  $u$  on  $M$  must be **constant**.*

In the special case where  $M = \mathbb{R}^n$ , the above result reduces to the classical **Liouville's Theorem** for **harmonic** functions on  $\mathbb{R}^n$ .

This theorem follows from the following *pointwise inequality* due to Cheng-Yau:

**Theorem 2.2** (Cheng-Yau 1975). *Let  $M$  be a complete manifold of dimension  $n \geq 2$  with  $\text{Ricci}(M) \geq 0$ . Suppose that  $u$  is any positive harmonic function in a geodesic ball  $B_R(x_0) \subset M$ . Then, the inequality*

$$\boxed{\text{eqn-dh1}} \quad (2.1) \quad \frac{|\nabla u|}{u} \leq \frac{C_n}{R}$$

holds in  $B_R(x_0)$ , where  $C_n$  depends only on the dimension  $n$ .

*Remark 2.3.* In the case where  $M^n = \mathbb{R}^n$  this resembles the derivative estimates for Harmonic functions, namely that for any harmonic function on  $\mathbb{R}^n$ , we have

$$\sup_{B_{\frac{R}{2}}(x_0)} |Du| \leq \frac{C_n}{R} \sup_{B_R(x_0)} |u|$$

for any  $x_0 \in \mathbb{R}^n$  and  $R > 0$ . The significance of the Cheng-Yau estimate is that it is pointwise, but on the other hand it assumes that  $u > 0$ .

Motivated by the elliptic result one may ask:

**Question:** Does the analogue of Yau's theorem hold for **positive** solutions of the heat equation

$$u_t = \Delta u \quad \text{on } M^n?$$

**Answer:** No. Example  $u(x, t) = e^{x_1+t}$ ,  $x = (x_1, \dots, x_n)$  on  $M^n := \mathbb{R}^n$ .

However **Souplet-Zhang** in [3] showed the following Liouville's theorem for harmonic functions on complete manifolds. We refer to the references in [3] for related work, and in particular with regard to the differentiable Harnack inequalities which will be discussed below.

$\boxed{\text{thm-lhe}}$  **Theorem 2.4.** *Let  $M^n$  be a complete non-compact Riemannian manifold of dimension  $n \geq 2$  with  $\text{Ricci}(M^n) \geq 0$ .*

- (1) *If  $u$  be a **positive ancient** solution to the heat equation on  $M^n \times (-\infty, T)$  such that*

$$u(p, t) = e^{o(d(p) + \sqrt{|t|})} \quad \text{as } d(p) \rightarrow \infty$$

*then  $u$  is a **constant**.*

- (2) *If  $u$  be an ancient solution to the heat equation on  $M^n \times (-\infty, T)$  such that*

$$u(p, t) = o(d(p) + \sqrt{|t|}) \quad \text{as } d(p) \rightarrow \infty$$

*then  $u$  is a **constant**.*

*Remark 2.5.* The example  $u(x, t) = e^{x_1+t}$  on  $\mathbb{R}^n$  shows that this theorem is sharp !

### 3. LI-YAU TYPE INEQUALITIES

The parabolic analogue of the Chen-Yau estimate for solutions to the heat equation on complete non-compact manifolds was shown by [Peter Li](#) and [S.T. Yau](#) [2]. Another improved Harnack was later shown by [R. Hamilton](#) [1] for the compact case. Both of these inequalities have played a fundamental role to the development of geometric flows. [Souplet-Zhang](#) in [3] **localized** the estimate by [R. Hamilton](#) on **complete non-compact manifolds**. In fact he proved the following result. Here and in the following we denote by  $Q_{R,T} := B_R(x_0) \times [t_0 - T, t_0]$  and  $Q_R := Q_{R,R^2}$ .

thm-SZ1

**Theorem 3.1.** *Let  $M$  be a complete non-compact Riemannian manifold of dimension  $n \geq 2$  with  $\text{Ricci}(M) \geq 0$ . Suppose that  $u$  is any positive solution in  $Q_{R,T}$  and that  $u \leq L$  in  $Q_{R,T}$ . Then,*

eqn-dhhe3

$$(3.1) \quad \frac{|\nabla u|}{u} \leq C_n \left( \frac{1}{R} + \frac{1}{\sqrt{T}} \right) \left( 1 + \ln \frac{L}{u} \right)$$

holds on  $Q_{R/2, T/2}$ , where  $C_n$  depends only on the dimension  $n$ .

*Remark 3.2.* Note that (3.1) is **scaling invariant** under parabolic scaling  $\beta u(ax, a^2t)$ .

We will see next that can be directly used to give us the proof of the Liouville type theorem for the heat equation, Theorem (2.4).

*Sketch of proof.* The proof goes by the maximum principle. Suppose that  $u \leq L$  in  $Q_{R,T}$ . Since our estimate is scaling invariant we may assume that  $L = 1$ , i.e.  $0 < u \leq 1$ . We set

$$f := \ln u, \quad w := |\nabla(1 - f)|^2 = \frac{|\nabla f|^2}{(1 - f)^2}.$$

Then, the above estimate (after squaring and setting  $L = 1$ ) becomes equivalent to

$$w \leq C_n \left( \frac{1}{R^2} + \frac{1}{T} \right)$$

which they show by the maximum principle !!

To this end, as usual we compute the evolution of  $w$  and after many standard calculations you show that

$$w_t - \Delta w \leq -\frac{2f}{1-f} \nabla f \cdot \nabla w - 2(1-f)w^2.$$

The details of this calculation in the simpler case where  $M = \mathbb{R}^n$  will be given by [Robin](#) during the discussion session !!

□

*Remark 3.3.* The assumption that  $u > 0$  is essential for this estimate to hold. Also, we will see that in the proof one works with  $f := \ln u$  and the positivity of  $u$  is heavily used.

*Remark 3.4.* In the estimate differential Harnack inequality (3.1) above as well as in the Cheng-Yau estimate the assumption  $\text{Ricci}(M) \geq 0$  can be generalized to the case  $\text{Ricci}(M) \geq -k$ , for some  $k > 0$ , with a slight modification of the statement. However, the assumption  $\text{Ricci}(M) \geq 0$  is necessary for the Liouville theorems to hold both in the elliptic and parabolic cases !!

#### 4. PROOF OF THEOREM 2.4

Lets assume that  $M = \mathbb{R}^n$ , although the proof doesn't change at all if  $M$  is any complete, noncompact manifold with nonnegative Ricci curvature.

(1) Assume that  $u$  be a **positive ancient** solution to the heat equation on  $\mathbb{R}^n \times (-\infty, T)$  such that

$$u(x, t) = e^{o(|x| + \sqrt{|t|})} \quad \text{as } |x| \rightarrow \infty.$$

We want to see then  $u$  is a **constant**. To this end, we will apply (3.1) to  $u + 1$  instead of  $u$ . Fix a point  $x_0 \in \mathbb{R}^n$  and a time  $t_0$  and consider the cube  $Q_R := B_R(x_0) \times [t_0 - R^2, t_0]$ . Then, by our assumption we have

$$0 \leq \ln(u + 1) \leq o(|x| + \sqrt{|t|}) = o(R), \quad \text{on } Q_R, R \gg 1.$$

Hence,  $u + 1 \leq e^{o(R)} := L_R$  and by (3.1), we have

$$\frac{|\nabla u(x_0, t_0)|}{u(x_0, t_0) + 1} \leq \frac{C_n}{R} \left(1 + \ln \frac{L_R}{u + 1}\right) \leq \frac{C(n, L)}{R} (1 + o(R))$$

where we have used that

$$\ln \frac{L_R}{u + 1} = \ln L_R - \ln(u + 1) \leq \ln L_R = o(R).$$

Letting,  $R \rightarrow +\infty$  we conclude that  $\nabla u(0, t_0) = 0$ , for all points  $(x_0, t_0)$ , i.e.  $u \equiv C$ .

(2) Assume now that  $u$  is an ancient solution to the heat equation on  $\mathbb{R}^n \times (-\infty, T)$  such that

$$u(p, t) = o(|x| + \sqrt{|t|}) \quad \text{as } |x| \rightarrow \infty.$$

Fix a point  $x_0 \in \mathbb{R}^n$  and a time  $t_0$  and consider and let  $A_R := \sup_{Q_R} |u|$ , where  $Q_R := B_R(x_0) \times [t_0 - R^2, t_0]$  as before. Our assumption implies that

$$A_R = o(R), \quad \text{as } R \rightarrow +\infty.$$

Set

$$U := u + 2A_R.$$

Then,

$$A_R \leq U \leq 3A_R, \quad \text{on } Q_R.$$

Apply (3.1) on  $Q_R$  and use that  $U \geq A_R$  and  $L_R = 3A_R$  to obtain

$$\frac{|\nabla U(x_0, t_0)|}{U(x_0, t_0)} \leq \frac{C}{R} \left(1 + \ln \frac{L_R}{U(x_0, t_0)}\right) \leq \frac{C}{R} (1 + \ln 3) \leq \frac{3C}{R}.$$

Here we have used that

$$\frac{U(x_0, t_0)}{R} \leq \frac{3A_R}{A_R} \leq 3 \quad \text{and} \quad \ln 3 \leq 3.$$

Thus,

$$\frac{|\nabla u(x_0, t_0)|}{U(x_0, t_0)} \leq \frac{3C}{R}$$

implying that

$$|\nabla u(x_0, t_0)| \leq 3C \frac{U(x_0, t_0)}{R} \leq 3C \frac{3A_R}{R} = \frac{o(R)}{R} \rightarrow 0, \quad \text{as } R \rightarrow +\infty.$$

We conclude that

$$\nabla u \equiv 0$$

i.e.  $u \equiv \text{constant}$ .

## REFERENCES

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