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# A $C^{0}$ estimate for nearly umbilical surfaces 

Received: 20 April 2005 / Accepted: 20 August 2005 / Published online: 24 March 2006
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#### Abstract

Let $\Sigma \subset \mathbf{R}^{3}$ be a smooth compact connected surface without boundary. Denote by $A$ its second fundamental form and by $\AA$ the tensor $A-(\operatorname{tr} A / 2)$ Id. In [4] we proved that, if $\|\AA\|_{L^{2}(\Sigma)}$ is small, then $\Sigma$ is $W^{2,2}$-close to a round sphere. In this note we show that, in addition, the metric of $\Sigma$ is $C^{0}$-close to the standard metric of $\mathbf{S}^{2}$.


## 1 Introduction

Let $\Sigma \subset \mathbf{R}^{3}$ be a smooth surface. A point $p$ of $\Sigma$ is called umbilical if the principal curvatures of $\Sigma$ at $p$ are equal and the surface $\Sigma$ is called umbilical if every point $x \in \Sigma$ is umbilical. A classical theorem in differential geometry states that if $\Sigma$ is a compact connected umbilical surface without boundary, then $\Sigma$ is a a round sphere. In [2] we proved the following quantitative version. Here:

- Id denotes the identity $(1,1)$-tensor and the $(0,2)$-tensor naturally associated to it;
$-\AA$ denotes the traceless part of $A$, i.e. the tensor $A-\frac{\operatorname{tr} A}{2} \mathrm{Id}$;
- id : $\mathbf{S}^{2} \subset \mathbf{R}^{3} \rightarrow \mathbf{R}^{3}$ is the standard isometric embedding of the round sphere.

Theorem 1 Let $\Sigma \subset \mathbf{R}^{3}$ denote a smooth compact connected surface without boundary and for convenience normalize the area of $\Sigma$ by $\operatorname{ar}(\Sigma)=4 \pi$. Then

$$
\begin{equation*}
\|A-\mathrm{Id}\|_{L^{2}(\Sigma)} \leq C\left\|\AA^{\AA}\right\|_{L^{2}(\Sigma)} \tag{1}
\end{equation*}
$$

The first author was supported by a grant of the Swiss National Science Foundation.
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where $C$ is a universal constant. If in addition $\left\|A^{\circ}\right\|_{L^{2}(\Sigma)}^{2} \leq 4 \pi$, then there exists a conformal parameterization $\psi: \mathbf{S}^{2} \rightarrow \Sigma$ and a vector $c_{\Sigma} \in \mathbf{R}^{3}$ such that

$$
\begin{equation*}
\left\|\psi-\left(c_{\Sigma}+\mathrm{id}\right)\right\|_{W^{2,2}\left(\mathbf{S}^{2}\right)} \leq C\|\AA\|_{L^{2}(\Sigma)} \tag{2}
\end{equation*}
$$

Since $\psi$ conformal, if we denote by $g$ the metric of $\Sigma$ and by $\sigma$ the standard metric on $\mathbf{S}^{2}$, then $\psi_{\#} g=h^{2} \sigma$ for some positive function $h$. Hence Theorem 1 gives

$$
\begin{equation*}
\|h-1\|_{W^{1,2}\left(\mathbf{R}^{2}\right)} \leq C\left\|A^{\circ}\right\|_{L^{2}\left(\mathbf{S}^{2}\right)} \tag{3}
\end{equation*}
$$

Therefore, by Sobolev embeddings, for every $p<\infty$ there exists a constant $C_{p}$ such that

$$
\|h-1\|_{L^{p}\left(\mathbf{S}^{2}\right)} \leq C_{p}\left\|\AA^{\circ}\right\|_{L^{2}\left(\mathbf{S}^{2}\right)}
$$

From (3) we cannot get a similar estimate for $\|h-1\|_{L^{\infty}}$. Nonetheless in this paper we show that such an estimate holds.

Theorem 2 There exists a universal constant $C$ with the following property. Let $\Sigma$ be any given compact connected surface of $\mathbf{R}^{3}$ without boundary, such that $\operatorname{ar}(\Sigma)=4 \pi$ and $\|\AA\|_{L^{2}(\Sigma)} \leq 8 \pi$. Then the conformal parameterization $\psi$ of Theorem 1 enjoys the bound

$$
\begin{equation*}
\|h-1\|_{C^{0}\left(\mathbf{S}^{2}\right)} \leq C\left\|A^{\circ}\right\|_{L^{2}(\Sigma)} \tag{4}
\end{equation*}
$$

We prove this estimate by suitably modifying techniques and ideas from [4]. There the authors showed bounds for $\|h\|_{\infty}$ when $A \in L^{2}$, by proving suitable bounds for $\operatorname{det} A$ in the Hardy space $\mathcal{H}^{1}$. These Hardy bounds were achieved through the $\mathbf{R}^{n}$ theory of [1] after locally lifting the Gauss map $N: \Sigma \rightarrow \mathbf{S}^{2}$ to a suitable map $M: \Sigma \rightarrow \mathbf{S}^{5}$. The same strategy can be implemented using $\mathbf{S}^{3}$-liftings. The core of Theorem 2 consists in showing that when $\left\|A^{\circ}\right\|_{L^{2}}$ is small, these liftings can be chosen $W^{1,2}$-close to suitable liftings of the identity map.

Estimate (4) is crucial to conclude that some geometric constants of $\Sigma$ are close to the corresponding ones of $\mathbf{S}^{2}$. For instance it implies that the spectrum of the Laplace-Beltrami operator of $\Sigma$ is close to that of $\mathbf{S}^{2}$. More precisely, given a compact surface $\Gamma$ without boundary, we denote by $\lambda_{i}(\Gamma)$ the $i$-th eigenvalue of the Laplace-Beltrami operator, with the following conventions: $\lambda_{0}(\Gamma)=0$ and if $a$ is an eigenvalue with multiplicity $n$, then it appears $n$ times in the sequence $\left\{\lambda_{i}(\Gamma)\right\}\left(\right.$ e.g. $\left.\lambda_{1}\left(\mathbf{S}^{2}\right)=\lambda_{2}\left(\mathbf{S}^{2}\right)=\lambda_{3}\left(\mathbf{S}^{2}\right)=2\right)$.

Corollary 1 For each $i$ there exists a constant $C_{i}$ with the following property. Let $\Sigma$ be any given compact connected surface of $\mathbf{R}^{3}$ without boundary, such that $\operatorname{ar}(\Sigma)=4 \pi$ and $\|\AA\|_{L^{2}(\Sigma)} \leq 4 \pi$. Then

$$
\begin{equation*}
\left|\lambda_{i}(\Sigma)-\lambda_{i}\left(\mathbf{S}^{2}\right)\right| \leq C_{i}\|\AA\|_{L^{2}(\Sigma)} \tag{5}
\end{equation*}
$$

## 2 Hardy bounds

We denote by

- $N$ the Gauss map on $\Sigma$;
$-M$ the $\operatorname{map} M:=N \circ \psi$;
- $K_{\Sigma}$ the Gauss curvature $\operatorname{det} d N$;
- $K$ the function $K:=K_{\Sigma} \circ \psi$;
$-\omega$ the standard volume form on $\mathbf{S}^{2}$.
In order to simplify the notation, for every 2-form $\alpha$ on $\mathbf{S}^{2}$ and every function space $H$, we denote by $\|\alpha\|_{H}$ the number $\|f\|_{H}$, where $f \omega=\alpha$.

Then Theorem 2 follows from the following Hardy bound.
Proposition 1 There exist positive constants $C$ and $\varepsilon$ such that the following holds. If $M: \mathbf{S}^{2} \rightarrow \mathbf{S}^{2}$ is a map such that $\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \leq \varepsilon$, then

$$
\begin{equation*}
\left\|M^{*} \omega-\omega\right\|_{\mathcal{H}^{1}\left(\mathbf{S}^{2}\right)} \leq C\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \tag{6}
\end{equation*}
$$

Proof (Proof of Theorem 2) Since $h$ is a positive function there exists a unique function $u$ such that $h=e^{u}$. Set

$$
\begin{equation*}
\delta:=\|\AA\|_{L^{2}(\Sigma)} \tag{7}
\end{equation*}
$$

From Proposition 3.2 of [2] we have that, under the assumptions of Theorem 2, there exists a universal constant $C_{1}$ such that

$$
\begin{equation*}
\|u\|_{C^{0}}+\|u\|_{W^{2,1}} \leq C_{1} . \tag{8}
\end{equation*}
$$

Thus it suffices to prove the existence of positive constants $\eta$ and $C_{2}$ such that

$$
\begin{equation*}
\|u\|_{C^{0}} \leq C_{2} \delta \quad \text { whenever } \delta<\eta \tag{9}
\end{equation*}
$$

Thanks to Theorem 1 and to the bounds (8), there exists a universal constant $C_{3}$ such that

$$
\begin{equation*}
\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \leq C_{3} \delta \tag{10}
\end{equation*}
$$

Let $\varepsilon$ be the constant of Proposition 1 and $\delta<\eta=\varepsilon / C_{3}$. Then we have

$$
\begin{equation*}
\left\|M^{*} \omega-\omega\right\|_{\mathcal{H}^{1}\left(\mathbf{S}^{2}\right)} \leq C_{4}\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \leq C_{5} \delta . \tag{11}
\end{equation*}
$$

Note that $K e^{2 u} \omega=M^{*} \omega$ and hence (11) gives

$$
\begin{equation*}
\left\|K e^{2 u}-1\right\|_{\mathcal{H}^{1}\left(\mathbf{S}^{2}\right)} \leq C_{5} \delta . \tag{12}
\end{equation*}
$$

Recall that $u$ satisfies

$$
\begin{equation*}
-\Delta_{\mathbf{S}^{2}} u=K e^{2 u}-1 \tag{13}
\end{equation*}
$$

Since the only harmonic functions on $\mathbf{S}^{2}$ are the constants, the bound (12) and the results of [3] imply that

$$
\begin{equation*}
\|u-c\|_{C^{0}} \leq C_{6} \delta \quad \text { for some constant } c \tag{14}
\end{equation*}
$$

The conformality of $\psi$ gives $4 \pi=\operatorname{ar}(\Sigma)=\int_{\mathbf{S}^{2}} e^{2 u}$ and (8) implies

$$
\left|e^{2(u-c)}-1\right| \leq C_{7}|u-c|
$$

for some constant $C_{7}$. Therefore we have

$$
4 \pi\left|e^{2 c}-1\right|=e^{2 c}\left|\int_{\mathbf{S}^{2}}\left(e^{2(u-c)}-1\right)\right| \leq C_{7} C_{6} 4 \pi \delta
$$

Hence there is a constant $C_{8}$ such that $|c| \leq C_{8} \delta$. From this and (14) we get (9).

The Hardy bound of Proposition 1 is proved by "locally" lifting the maps $M$ and id to maps into $\mathbf{S}^{3}$ via the Hopf fibration $\pi: \mathbf{S}^{2} \rightarrow \mathbf{S}^{3}$. The reason why we cannot argue globally is that there is no such smooth lifting for the identity. Let $p \in \mathbf{S}^{2}$ and denote by $D_{\pi / 2+1}(p)$ the geodesic disk of $\mathbf{S}^{2}$ with center $p$ and radius $\pi / 2+1$. Then in the next two sections we will prove the following proposition.

Proposition 2 (Hardy bound) Let $\Psi \in C^{\infty}\left(\mathbf{S}^{2}, \mathbf{S}^{2}\right)$ be a fixed map with $\Psi\left(\mathbf{S}^{2}\right) \subset$ $D_{\pi / 2+1}(p)$. There exist positive constants $C$ and $\varepsilon$, depending only on $\|\Psi\|_{C^{2}}$, such that:
(HB) If $M \in C^{\infty}\left(\mathbf{S}^{2}, \mathbf{S}^{2}\right)$ satisfies $\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \leq \varepsilon$, then

$$
\begin{equation*}
\left\|M^{*} \omega-\Psi^{*} \omega\right\|_{\mathcal{H}^{1}\left(\mathbf{S}^{2}\right)} \leq C\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} . \tag{15}
\end{equation*}
$$

Note that, since $\Psi\left(\mathbf{S}^{2}\right) \subset D_{\pi / 2+1}(p)$, there exists a smooth lifting of $\Psi$ through the Hopf fibration (see Proposition 3). This lifting exists under the weaker assumption $\int_{\mathbf{S}^{2}} \Psi^{*} \omega=0$. However, the stronger assumption $\Psi\left(\mathbf{S}^{2}\right) \subset D_{\pi / 2+1}(p)$ will be crucial later in order to prove some estimates on the lifting (compare with the Second Step of the proof of Lemma 1).

From Proposition 2 one concludes Proposition 1 with a "cut and paste" procedure.

Proof (Proof of Proposition 1) First of all we introduce some notation. We let $p$ be any point of $\mathbf{S}^{2} \subset \mathbf{R}^{3}$. Then we let

$$
D:=D_{\pi / 2+1 / 2}(p) \quad \tilde{D}:=D_{\pi / 2+1}(p)
$$

We claim that if $M$ is a smooth map and $\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}$ is sufficiently small, then there exist two maps $M^{\prime}, \Psi^{\prime}: \mathbf{S}^{2} \rightarrow \mathbf{S}^{2}$ such that:
$-M^{\prime}=M$ and $\Psi^{\prime}=\mathrm{id}$ on $D$;
$-\Psi^{\prime}\left(\mathbf{S}^{2}\right) \subset \tilde{D}$;

- The following estimates hold for some universal constant $C$ :

$$
\begin{equation*}
\left\|M^{\prime}-\Psi^{\prime}\right\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \leq C\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \quad\left\|\Psi^{\prime}\right\|_{C^{2}} \leq C \tag{16}
\end{equation*}
$$

This fact, combined with Proposition 2, yields the the existence of two positive constants $C$ and $\varepsilon$ such that

$$
\begin{equation*}
\left\|M^{*} \omega-\omega\right\|_{\mathcal{H}^{1}\left(D_{\pi / 2+1 / 2}(p)\right)} \leq C\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \tag{17}
\end{equation*}
$$

for all $p \in \mathbf{S}^{2}$ and all $M$ with $\|M-\mathrm{id}\|_{W^{1,2}}<\varepsilon$. Note that if $p$ and $q$ are two antipodal points, then

$$
D_{\pi / 2+1 / 2}(p) \cup D_{\pi / 2+1 / 2}(q)=\mathbf{S}^{2}
$$

Therefore from (17) we would get

$$
\begin{equation*}
\left\|M^{*} \omega-\omega\right\|_{\mathcal{H}^{1}\left(\mathbf{S}^{2}\right)} \leq C^{\prime}\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \tag{18}
\end{equation*}
$$

which is the desired conclusion. It remains to prove the existence of the maps $M^{\prime}$ and $\Psi^{\prime}$.

First Step By Fubini's Theorem, there exists a universal constant $C$ with the following property: There exists $\rho \in[\pi / 2+1 / 2, \pi / 2+3 / 4]$ such that

$$
\begin{align*}
\|M-\mathrm{id}\|_{W^{1,2}\left(\partial D_{\rho}(p)\right)} & \leq\|M-\mathrm{id}\|_{L^{2}\left(\partial D_{\rho}(p)\right)}+\|D(M-\mathrm{id})\|_{L^{2}\left(\partial D_{\rho}(p)\right)} \\
& \leq C\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \tag{19}
\end{align*}
$$

Now let us fix radial coordinates $\theta, r$ on $\tilde{D}$. We define $\tilde{M}, \tilde{\Psi}:=\tilde{D} \rightarrow \mathbf{S}^{2}$ as

$$
\begin{aligned}
& \tilde{M}(\theta, r)= \begin{cases}M(\theta, r) & \text { if } r<\rho \\
M(\theta, \pi / 2+3 / 4) & \text { if } r \geq \rho\end{cases} \\
& \tilde{\Psi}(\theta, r)= \begin{cases}\operatorname{id}(\theta, r) & \text { if } r<\rho \\
\operatorname{id}(\theta, \pi / 2+3 / 4) & \text { if } r \geq \rho\end{cases}
\end{aligned}
$$

Clearly $\|\tilde{M}-\tilde{\Psi}\|_{W^{1,2}(\tilde{D})} \leq C\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}$ for some universal constant $C$.
Second Step We claim the existence of positive constants $\varepsilon$ and $\eta$ with the following property. If $\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \leq \varepsilon$, then there exists a point $q \in \mathbf{S}^{2} \backslash \tilde{D}$ such that

$$
\operatorname{dist}(q, \tilde{M}(\partial \tilde{D}))+\operatorname{dist}(q, \partial \tilde{D}) \geq \eta
$$

This claim will be proved later. Assuming it, we set $\zeta:=\min \{1 / 8, \eta / 2\}$. Using such a point $q$ we can construct a $C^{2}$ map

$$
R:[\pi / 2+3 / 4, \pi / 2+1] \times\left\{\mathbf{S}^{2} \backslash D_{\zeta}(q)\right\} \rightarrow \mathbf{S}^{2}
$$

such that:

- $R(t, \cdot)$ maps $\tilde{D}$ into $\tilde{D}$ for every $t$;
$-R(\pi / 2+1, \cdot)$ maps $\mathbf{S}^{2} \backslash D_{\zeta}(q)$ onto $p$;
- $\|R\|_{C^{2}}$ is bounded by a universal constant depending only on $\zeta$.

Given such an $R$ we define the maps $M^{\prime}, \Psi^{\prime}: \tilde{D} \rightarrow \mathbf{S}^{2}$ as

$$
\begin{aligned}
& M^{\prime}(\theta, r)= \begin{cases}\tilde{M}(\theta, r) & \text { if } r<\pi+3 / 4 \\
R(r, \tilde{M}(\theta, \pi / 2+3 / 4)) & \text { if } r \geq \pi+3 / 4\end{cases} \\
& \Psi^{\prime}(\theta, r)= \begin{cases}\tilde{\Psi}(\theta, r) & \text { if } r<\pi+3 / 4 \\
R(r, \tilde{\Psi}(\theta, \pi / 2+3 / 4)) & \text { if } r \geq \pi+3 / 4\end{cases}
\end{aligned}
$$

Finally, we extend both $\Psi^{\prime}$ and $M^{\prime}$ to $\mathbf{S}^{2}$ by setting $\Psi^{\prime}=M^{\prime}=p$ on $\mathbf{S}^{2} \backslash \tilde{D}$. Then $M^{\prime}$ and $\Psi^{\prime}$ would satisfy all the requirements of the Lemma. Therefore, in order to conclude the proof it suffices to show the existence of the point $q$.
Third Step For any regular value $\tilde{\sim} \underset{\sim}{\mathcal{D}} \in \mathbf{S}^{2} \backslash \tilde{M}(\partial \tilde{D})$ we define the degree $\operatorname{deg}(\tilde{q}, \tilde{M}, \tilde{D})$ in the usual way. It is a classical fact that deg is constant in the connected component of $\mathbf{S}^{2} \backslash \tilde{M}(\partial \tilde{D})$. Hence we extend it to $\mathbf{S}^{2} \backslash \tilde{M}(\partial \tilde{D})$ by continuity and we set

$$
U_{0}:=\left\{\tilde{q} \in \mathbf{S}^{2}: \operatorname{deg}(\tilde{q}, \tilde{M}, \tilde{D})=0\right\}
$$

It turns out that $U_{0}$ is an open set with boundary contained in the curve

$$
\gamma=\tilde{M}(\partial \tilde{D})=M(\partial \tilde{D})
$$

By (19) the length of $\gamma$ is less than $C+C\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}$, for some universal constant $C$.

Consider the open set $U:=U_{0} \backslash \tilde{D}$. Clearly, $\left\{\mathbf{S}^{2} \backslash[\tilde{M}(\tilde{D}) \cup \tilde{D}]\right\} \subset U$. Moreover, by construction we have $\tilde{M}(\tilde{D}) \subset M(\tilde{D})$. From the area formula it follows that

$$
\operatorname{ar}(M(\tilde{D}) \backslash \tilde{D}) \leq C\|M-\mathrm{id}\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}
$$

Therefore, there exist positive universal constants $C_{1}, C_{2}, C_{3}$ such that, if $\| M-$ $\mathrm{id} \|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \leq C_{1}$ then $U$ is an open set with the following properties:

- $\partial \tilde{D}_{\tilde{D}}$ is contained in the union of two connected curves $\gamma=\tilde{M}(\partial \tilde{D})$ and $\tilde{\gamma}=$ д $\tilde{D}$;
$-\operatorname{ar}(U) \geq C_{2}$ and len $(\gamma)+\operatorname{len}(\tilde{\gamma}) \leq C_{3}$.
An elementary argument shows the existence of a positive constant $\eta$ such that every $U$ satisfying the conditions above contains a disk of radius $\eta$ (see for instance Lemma C. 1 of [2]). The center of this disk is the desired point $q$.


## 3 Liftings through Hopf fibration

Denote by $\pi: \mathbf{S}^{3} \subset \mathbb{C}^{2} \rightarrow \mathbf{S}^{2}$ the Hopf fibration. Note that if we choose $\varepsilon$ small enough in Proposition 2, then we have

$$
\begin{equation*}
\left|\int_{\mathbf{S}^{2}}\left(M^{*} \omega-\Psi^{*} \omega\right)\right|<1 \tag{20}
\end{equation*}
$$

From classical topological arguments we know that $\int_{\mathbf{S}^{2}} M^{*} \omega$ is an integer and that $\int_{\mathbf{S}^{2}} \Psi^{*} \omega=0$ (this last equality follows from the assumption $\left.\Psi\left(\mathbf{S}^{2}\right) \subset \tilde{D}\right)$. Therefore $\int_{\mathbf{S}^{2}} M^{*} \omega=0$.

The condition $\int_{\mathbf{S}^{2}} \Psi^{*} \omega=\int_{\mathbf{S}^{2}} M^{*} \omega=0$ implies that the maps $\Psi$ and $M$ are homotopically trivial. Therefore there exist smooth maps $\Phi, F: \mathbf{S}^{2} \rightarrow \mathbf{S}^{3}$ with $\pi \circ \Phi=\Psi$ and $\pi \circ F=M$. One main idea of [4] is that one can prove an Hardy bound $\left\|M^{*} \omega\right\|_{\mathcal{H}^{1}}$ by showing that the lifting $\Psi$ can be chosen with bounded $W^{1,2}$ norm. (In passing we remark that in the paper [4] the authors used liftings to $\mathbf{S}^{5}$; however this is only a technical difference, mainly due to the fact that in [4] this technique is applied to the case of 2-dimensional surfaces in $\mathbf{R}^{n}$.) Therefore one naturally expects that, if the liftings $\Psi$ and $M$ can be chosen $W^{1,2}$-close, then one gets the bound (15).

Proposition 3 Let $\Psi$ and $M$ be as in Proposition 2. Then there exist two maps $\Phi, F: \mathbf{S}^{2} \rightarrow \mathbf{S}^{3}$ such that

- $\Psi=\pi \circ \Phi, M=\pi \circ F$;
- $\|\Phi\|_{C^{1}} \leq C,\|F-\Phi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \leq C\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}$.

The constant $C$ depends only on $\|\Psi\|_{C^{2}}$ and not on $M$.

Building on this proposition, the proof of Proposition 2 is a short argument. However we set first a bit of notation. We fix coordinates on $\mathbb{C}^{2}$ so that

$$
\begin{aligned}
& \mathbf{S}^{3}=\left\{\left.\left(z_{1}, z_{2}\right) \in \mathbb{C}^{2}| | z_{1}\right|^{2}+\left|z_{2}\right|^{2}=1\right\} \\
& \mathbf{S}^{2}=\left\{(z, t) \in \mathbb{C} \times\left.\mathbf{R}| | z\right|^{2}+t^{2}=1\right\}
\end{aligned}
$$

Then the Hopf fibration is given by $\pi\left(z_{1}, z_{2}\right)=\left(2 \bar{z}_{1} z_{2},\left|z_{1}\right|^{2}-\left|z_{2}\right|^{2}\right)$. Note that if $p=\left(z_{1}, z_{2}\right) \in \mathbf{S}^{3}$, then the fiber

$$
\begin{equation*}
F_{p}:=\left\{\left(w_{1}, w_{2}\right) \mid \pi\left(w_{1}, w_{2}\right)=\pi\left(z_{1}, z_{2}\right)\right\} \tag{21}
\end{equation*}
$$

is given by $\left\{\left(e^{i \theta} z_{1}, e^{i \theta} z_{2}\right), \theta \in \mathbf{R}\right\}$.
Proof (Proof of Proposition 2) Let $\Phi$ and $F$ be the liftings of Proposition 3. Using the coordinates above we write $\Phi=\left(\Phi_{1}, \Phi_{2}\right)$ and $F=\left(F_{1}, F_{2}\right)$. The following identities can be easily checked:

$$
\begin{align*}
& 2 \Psi^{*} \omega=2 \Phi^{*} \pi^{*} \omega=i\left(d \Phi_{1} \wedge d \bar{\Phi}_{1}+d \Phi_{2} \wedge d \bar{\Phi}_{2}\right) \\
& 2 M^{*} \omega=2 F^{*} \pi^{*} \omega=i\left(d F_{1} \wedge d \bar{F}_{1}+d F_{2} \wedge d \bar{F}_{2}\right) \tag{22}
\end{align*}
$$

Note that

$$
\begin{aligned}
2\left(\Psi^{*} \omega-M^{*} \omega\right)= & i\left\{d \Phi_{1} \wedge d\left(\bar{\Phi}_{1}-\bar{F}_{1}\right)+d\left(\Phi_{1}-F_{1}\right) \wedge d \bar{F}_{1}\right. \\
& \left.+d \Phi_{2} \wedge d\left(\bar{\Phi}_{2}-\bar{F}_{2}\right)+d\left(\Phi_{2}-F_{2}\right) \wedge d \bar{F}_{2}\right\}
\end{aligned}
$$

Hence, using the results of [1] we get

$$
\left\|\Psi^{*} \omega-M^{*} \omega\right\|_{\mathcal{H}^{1}\left(\mathbf{S}^{2}\right)} \leq C\left(\|d F\|_{L^{2}\left(\mathbf{S}^{2}\right)}+\|d \Phi\|_{L^{2}\left(\mathbf{S}^{2}\right)}\right)\|d F-d \Phi\|_{L^{2}\left(\mathbf{S}^{2}\right)}
$$

Therefore the bounds satisfied by $F$ and $\Phi$ yield the desired estimate.
The rest of the paper is devoted to prove the existence of the liftings claimed in Proposition 3. First we introduce a suitable norm on differentials of maps with target in $\mathbf{S}^{3}$, see (24). This norm is invariant under the action of $\mathbf{S}^{3}$ on itself as Lie group.

We recall that $\mathbb{C}^{2}$ can be identified to the field of quaternions $\mathbb{H}$. We denote by $\times$ the multiplication between quaternions and we recall that the usual norm $|\cdot|$ has the property that $|a \times b|=|a||b|$. Hence, $\times$ naturally induces a Lie group structure on $\mathbf{S}^{3}$ and the maps $l^{w}: \mathbf{S}^{3} \rightarrow \mathbf{S}^{3}$ given by $l^{w}(a)=w \times a$ are isometries of $\mathbf{S}^{3}$. The same holds for the maps $r^{w}: \mathbf{S}^{3} \rightarrow \mathbf{S}^{3}$ given by $r^{w}(a)=a \times w$.

Definition 1 Given $a, b \in \mathbf{S}^{3}$ and $\xi \in T_{a} \mathbf{S}^{3}$ we denote by $b \xi$ the vector of $T_{b \times a} \mathbf{S}^{3}$ given by $\left.d l^{b}\right|_{a}(\xi)$. In a similar way we define $\xi b$ as $\left.d r^{b}\right|_{a}(\xi) \in T_{a \times b} \mathbf{S}^{3}$.

The diffeomorphisms $l^{x}$ allow to define an "intrinsic" notion of distance between vectors belonging to $T_{a} \mathbf{S}^{3}$ and $T_{b} \mathbf{S}^{3}$. This allows a natural way to compare the differential of two distinct maps with target in $\mathbf{S}^{3}$.

Definition 2 Given $\xi \in T_{b} \mathbf{S}^{3}, \zeta \in T_{a} \mathbf{S}^{3}$ we denote by $|\xi-\zeta|_{\mathcal{L}}$ the nonnegative real number

$$
\left|a^{-1} \xi-b^{-1} \eta\right|=\left|\left(b \times a^{-1}\right) \xi-\eta\right|=\left|\xi-\left(a \times b^{-1}\right) \zeta\right|
$$

where, for vectors $\lambda, \mu \in T_{p} \mathbf{S}^{3},|\lambda-\mu|$ denotes the usual Hilbert norm (that is, the norm induced by the Riemann structure of $\mathbf{S}^{3}$ as submanifold of $\mathbf{R}^{4}$ ).

Given a riemannian manifold $\Omega$ and smooth maps $F, \Phi: \Omega \rightarrow \mathbf{S}^{3}$, we define

$$
\begin{align*}
|d F|_{p}-\left.\left.d \Phi\right|_{p}\right|_{\mathcal{L}} & :=\sup _{|\xi|=1}|d F|_{p}(\xi)-\left.\left.d \Phi\right|_{p}(\xi)\right|_{\mathcal{L}}  \tag{23}\\
\left\|\left.\|d F-d \Phi\|\right|_{L^{2}(\Omega)}\right. & :=\left(\int_{\Omega}|d F-d \Phi|_{\mathcal{L}}^{2}\right)^{1 / 2} \tag{24}
\end{align*}
$$

The proof of Proposition 3 is based on two lemmas. The first one, Lemma 1, shows the existence of liftings for which one can estimate the norm $\left\|\|d F-d \Phi\|_{L^{2}\left(D_{r}\right)}\right.$ as in (26). The second, Lemma 2, is a Poincare' type inequality. With the help of this inequality, one can absorb the second term of (26), provided $r$ is smaller than a universal constant. This gives an estimate of the form

$$
\begin{equation*}
\|\|d F-d \Phi\|\|_{L^{2}\left(D_{r}\right)} \leq C\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \tag{25}
\end{equation*}
$$

The number of disks $D_{r}$ needed to cover $\mathbf{S}^{2}$ is smaller than a universal constant. Therefore we can bound $\|\|F-d \Phi\|\|_{L^{2}\left(\mathbf{S}^{2}\right)}$. We then use again Lemma 2 to show the existence of a new lifting $\tilde{F}$ such that

$$
\|d \tilde{F}-d \Phi\|\left\|_{L^{2}\left(\mathbf{S}^{2}\right)}+\right\| \tilde{F}-\Phi\left\|_{L^{2}} \leq C\right\| M-\Psi \|_{W^{1,2}\left(\mathbf{S}^{2}\right)}
$$

Finally it is not difficult to show that

$$
\|\tilde{F}-\Phi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \leq\|d \tilde{F}-d \Phi\|_{L^{2}\left(\mathbf{S}^{2}\right)}+\|F-\Phi\|_{L^{2}}
$$

Lemma 1 Let $M$ and $\Psi$ be as in Proposition 2 and choose $\varepsilon$ sufficiently small so that $M$ is homotopically trivial. Then there exists a universal constant $C$ and two maps $F, \Phi: \mathbf{S}^{2} \rightarrow \mathbf{S}^{3}$ such that:
$-\Psi=\pi \circ \Phi, M=\pi \circ F$ and $\|\Phi\|_{C^{1}} \leq C ;$

- For every disk $D_{r} \subset \mathbf{S}^{2}$ we have the estimate

$$
\begin{equation*}
\|d F-d \Phi\|_{L^{2}\left(D_{r}\right)} \leq C\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}+C \min _{\theta}\left\|F-e^{i \theta} \Phi\right\|_{L^{2}\left(D_{r}\right)} \tag{26}
\end{equation*}
$$

Lemma 2 Let $D_{r}$ be a geodesic disk of $\mathbf{S}^{3}$ and $\Phi, F: D_{r} \rightarrow \mathbf{S}^{3}$ two smooth maps. Then

$$
\begin{equation*}
\min _{w \in \mathbf{S}^{3}}\|F-w \times \Phi\|_{L^{2}\left(D_{r}\right)} \leq C r\|d F-d \Phi\| \|_{L^{2}\left(D_{r}\right)} \tag{27}
\end{equation*}
$$

for some universal constant $C$.
The proof of Lemma 1 is given in the next Section. Hereby we prove Lemma 2 and we show how to conclude Proposition 2.

Proof (Proof of Lemma 2) Let $G: D_{r} \rightarrow \mathbf{S}^{3} \subset \mathbb{H}$ be given by $G(p)=F(p) \times$ $\Phi(p)^{-1}$. Using the notation of Definition 1 we write

$$
d G_{p}(\xi)=\left(\left.d F\right|_{p}(\xi)\right) \Phi(p)^{-1}-\left[F(p) \Phi(p)^{-1}\right]\left(\left.d \Phi\right|_{p}(\xi)\right) \Phi(p)^{-1}
$$

Since the multiplication from the right is an isometry, we get $|\zeta b-\xi b|=|\zeta-\xi|$ for every $\xi \in T_{a} \mathbf{S}^{3}, \zeta \in T_{a} \mathbf{S}^{3}$. Hence

$$
\begin{equation*}
|d G|_{p}(\xi)\left|=|d F|_{p}(\xi)-\left[F(p) \Phi(p)^{-1}\right]\left(\left.d \Phi\right|_{p}(\xi)\right)\right| \tag{28}
\end{equation*}
$$

We remark that the right hand side of (28) is precisely the definition of $|d F|_{p}(\xi)-$ $\left.\left.d \Phi\right|_{p}(\xi)\right|_{\mathcal{L}}$. Thus,

$$
\|d G\|_{L^{2}\left(D_{r}\right)}=\|d F-d \Phi\| \|_{L^{2}\left(D_{r}\right)}
$$

Hence, by the usual Poincaré inequality on Euclidean spaces, there exists $w \in$ $\mathbb{H}=\mathbb{C}^{2}$ such that

$$
\|G-w\|_{L^{2}\left(D_{r}\right)} \leq C r\|d G\|_{L^{2}\left(D_{r}\right)}=C r\|d F-d \Phi\| \|_{L^{2}\left(D_{r}\right)}
$$

Note that

$$
\begin{align*}
\pi r^{2}|1-|w| \| & =\int_{D_{r}}\|G|-|w|| \leq\| G-w \|_{L^{1}\left(D_{r}\right)} \\
& \leq C_{1} r\|G-w\|_{L^{2}\left(D_{r}\right)} \leq C_{2} r^{2}\|d F-d \Phi\|_{L^{2}\left(D_{r}\right)} \tag{29}
\end{align*}
$$

Set $\tilde{w}:=w /|w|$. Then, by (29), we have $|\tilde{w}-w|=|1-|w|| \leq C_{3}| | \mid d F-$ $d \Phi \|_{L^{2}\left(D_{r}\right)}$. Hence
$\|G-\tilde{w}\|_{L^{2}\left(D_{r}\right)} \leq C_{4} r|\tilde{w}-w|+C_{5}\|w-G\|_{L^{2}\left(D_{r}\right)} \leq C_{6} r\|d F-d \Phi\|_{L^{2}\left(D_{r}\right)}$.
Since $\tilde{w} \in \mathbf{S}^{3}$, this gives the desired inequality.
Proof (Proof of Proposition 2) We start from the liftings $F$ and $\Phi$ provided by Lemma 1 and we break the proof into two steps.
First Step In this step we show that

$$
\begin{equation*}
\|d F-d \Phi\|\left\|_{L^{2}\left(D_{r}\right)} \leq C_{2}\right\| M-\Psi \|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \quad \text { if } r \leq C_{1} \tag{31}
\end{equation*}
$$

for some universal constant $C_{1}$. Since $\mathbf{S}^{2}$ is compact (31) implies

$$
\begin{equation*}
\|d F-d \Phi\|\left\|_{L^{2}\left(\mathbf{S}^{2}\right)} \leq C\right\| M-\Psi \|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \tag{32}
\end{equation*}
$$

Recall the Poincaré inequality proved in Lemma 2:

$$
\begin{equation*}
\min _{w \in \mathbf{S}^{3}}\|F-w \times \Phi\|_{L^{2}\left(D_{r}\right)} \leq C r\|d F-d \Phi\| \|_{L^{2}\left(D_{r}\right)} \tag{33}
\end{equation*}
$$

Let $w$ be a point where the minimum in the left hand side of (33) is attained and let $\theta_{0}$ be a point where $f(\theta)=\left|w-e^{i \theta}\right|$ attains its minimum. Recall that the quaternionic multiplication by an element of $\mathbf{S}^{3}$ is an isometry of $\mathbf{S}^{3}$. Thus, for every $a \in \mathbf{S}^{3}$, the function $f_{a}(\theta)=\left|w \times a-e^{i \theta} a\right|$ attains its minimum in $\theta_{0}$.

It is not difficult to check that

$$
\min _{\theta}\left|w \times a-e^{i \theta} a\right| \leq C_{1}|\pi(w \times a)-\pi(a)|,
$$

for some universal constant $C_{1}$. Moreover, recall that $\pi$ is Lipschitz and call $C_{2}$ its Lipschitz constant. Thus

$$
\begin{align*}
& \left\|w \times \Phi-e^{i \theta_{0}} \Phi\right\|_{L^{2}\left(D_{r}\right)} \leq C_{1}\|\pi(w \times \Phi)-\pi(\Phi)\|_{L^{2}\left(D_{r}\right)} \\
& \quad \leq C_{1}\|\pi(w \times \Phi)-\pi(F)\|_{L^{2}\left(D_{r}\right)}+C_{1}\|\pi(F)-\pi(\Phi)\|_{L^{2}\left(D_{r}\right)} \\
& \leq C_{1} C_{2}\|w \times \Phi-F\|_{L^{2}\left(D_{r}\right)}+C_{1}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} . \tag{34}
\end{align*}
$$

Combining (34) and (33) we get

$$
\begin{equation*}
\min _{\theta}\left\|F-e^{i \theta} \Phi\right\|_{L^{2}\left(D_{r}\right)} \leq C_{3} r\|d F-d \Phi\|_{L^{2}\left(D_{r}\right)}+C_{4}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} . \tag{35}
\end{equation*}
$$

Plugging (35) into (26) we get

$$
\begin{equation*}
\|\mid d F-d \Phi\|_{L^{2}\left(D_{r}\right)} \leq C_{5}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}+C_{6} r\|d F-d \Phi\|_{L^{2}\left(D_{r}\right)} . \tag{36}
\end{equation*}
$$

Thus it is sufficient to choose $r \leq\left(2 C_{6}\right)^{-1}$ to get

$$
\begin{equation*}
\|d F-d \Phi\|_{L^{2}\left(D_{r}\right)} \leq 2 C_{7}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} . \tag{37}
\end{equation*}
$$

This gives (31) and hence completes the proof of (32).
Second Step Conclusion
Let $\xi \in T_{a} \mathbf{S}^{3}, \zeta \in T_{b} \mathbf{S}^{3}$. The following elementary inequality holds:

$$
\begin{equation*}
|\xi-\zeta| \leq C|\xi||a-b|+C|\xi-\zeta|_{\mathcal{L}} . \tag{38}
\end{equation*}
$$

Indeed, since the map

$$
\mathbf{S}^{3} \times T \mathbf{S}^{3} \ni(w, a, \xi) \rightarrow w \xi \in T_{w \times a} \mathbf{S}^{3} \subset \mathbb{C}^{2}
$$

is Lipschitz on compact sets, we have

$$
\left|\xi-\left(b \times a^{-1}\right) \xi\right| \mathcal{L} \leq C\left|1-b \times a^{-1}\right|=C|a-b| \quad \text { for }|\xi| \leq 1 .
$$

Thus, if we define $\tilde{\xi}=\xi /|\xi|$ we get

$$
\begin{aligned}
|\xi-\zeta| & \leq\left|\left(b \times a^{-1}\right) \xi-\zeta\right|+\left|\left(b \times a^{-1}\right) \xi-\xi\right| \\
& =|\xi-\zeta| \mathcal{L}+|\xi|\left|\left(b \times a^{-1}\right) \tilde{\xi}-\tilde{\xi}\right| \\
& \leq|\xi-\zeta|_{\mathcal{L}}+C|\xi||b-a| .
\end{aligned}
$$

Let $\theta_{0}$ be a point where the expression

$$
g(\theta)=\left\|e^{i \theta} F-\Phi\right\|_{L^{2}\left(\mathbf{S}^{2}\right)}
$$

attains its minimum. Set $\tilde{F}=e^{i \theta_{0}} F$. Replacing $D_{r}$ with $\mathbf{S}^{2}$ in (35) we get

$$
\begin{align*}
\|\tilde{F}-\Phi\|_{L^{2}\left(\mathbf{S}^{2}\right)} & \leq C_{1}\|d d F-d \Phi\|_{L^{2}\left(\mathbf{S}^{2}\right)}+C_{1}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \\
& \leq C_{2}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} \tag{39}
\end{align*}
$$

From (38) we get

$$
|d \tilde{F}-d \Phi|^{2} \leq 2|d \tilde{F}-d \Phi|_{\mathcal{L}}^{2}+2 C^{2}|d \Phi|^{2}|\tilde{F}-\Phi|^{2}
$$

Integrating this inequality we get

$$
\begin{align*}
\|d \tilde{F}-d \Phi\|_{L^{2}\left(\mathbf{S}^{2}\right)}^{2} & \leq 2 C^{2} \int_{\mathbf{S}^{2}}|d \Phi|^{2}|\tilde{F}-\Phi|^{2}+\|\mid d F-d \Phi\|_{L^{2}\left(\mathbf{S}^{2}\right)}^{2} \\
& \stackrel{(32)}{\leq} C_{3}\|\Phi\|_{C^{1}}^{2}\|\tilde{F}-\Phi\|_{L^{2}}^{2}+C_{4}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}^{2} \\
& \stackrel{(39)}{\leq} C_{5}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}^{2} . \tag{40}
\end{align*}
$$

This concludes the proof.

## 4 Proof of Lemma 1

Recall the definition of $F_{p}$ given in (21) and note that the vector tangent to $F_{p}$ in $p=\left(z_{1}, z_{2}\right)$ is $\left(i z_{1}, i z_{2}\right)$. Thus, we decompose $T_{p} \mathbf{S}^{3}$ into two orthogonal subspaces:

$$
\begin{equation*}
T F_{p}=\left\{t\left(i z_{1}, i z_{2}\right) \mid t \in \mathbf{R}\right\} \quad T N_{p}=\left\{w \in T_{p} \mathbf{S}^{3} \mid w \cdot\left(i z_{1}, i z_{2}\right)=0\right\} \tag{41}
\end{equation*}
$$

where the hermitian product $\left(a_{1}, a_{2}\right) \cdot\left(b_{1}, b_{2}\right)$ is given by $\operatorname{Re}\left(a_{1} \bar{b}_{1}+a_{2} \bar{b}_{2}\right)$.
Definition 3 If $\Phi: \Omega \rightarrow \mathbf{S}^{3}$ is a smooth map, we we write $d \Phi=d_{1} \Phi+d_{2} \Phi$, where

- $\left.d_{1} \Phi\right|_{q}(\xi)$ is the projection of $\left.d \Phi\right|_{q}(\xi)$ on $T F_{\Phi(q)}$,
$-\left.d_{2} \Phi\right|_{q}(\xi)$ is the projection of $\left.d \Phi\right|_{q}(\xi)$ on $T N_{\Phi(q)}$.
Proof (Proof of Lemma 1)
First Step In this step we derive a preliminary estimate on $\||d F|-|d \Phi|\|_{L^{2}\left(\mathbf{S}^{2}\right)}$, provided $F$ and $\Phi$ are chosen in a suitable way.

First of all fix any pair of liftings $(F, \Phi)$. It can be easily checked that $\left|d_{2} F\right|=$ $|d M|$ and $\left|d_{2} \Phi\right|=|d \Psi|$. Moreover, if we define the 1 -form $\alpha:=-i \bar{z}_{1} d z_{2}-$ $i \bar{z}_{2} d z_{2}$, then we get

$$
\begin{equation*}
d_{1} F=\left(i F_{1}, i F_{2}\right) F^{*} \alpha \quad d_{1} \Phi=\left(i \Phi_{1}, i \Phi_{2}\right) \Phi^{*} \alpha \tag{42}
\end{equation*}
$$

Thus

$$
\int_{\mathbf{S}^{2}}| | d_{1} F\left|-\left|d_{1} \Phi \|^{2}=\int_{\mathbf{S}^{2}}\right| F^{*} \alpha-\Phi^{*} \alpha\right|^{2}
$$

We will show that the liftings $F$ and $\Phi$ can be chosen so that

$$
\int_{\mathbf{S}^{2}}\left|F^{*} \alpha-\Phi^{*} \alpha\right|^{2}=\left\|M^{*} \omega-\Psi^{*} \omega\right\|_{W^{-1,2}}^{2}
$$

Indeed, fix a lifting $\tilde{F}: \mathbf{S}^{2} \rightarrow \mathbf{S}^{3}$ of $M$ and set $\beta=\tilde{F}^{*} \alpha$. We can use the standard Hodge decomposition to write

$$
\beta=d \theta+* d \psi
$$

where $\theta$ and $\psi$ are smooth functions on $\mathbf{S}^{2}$. If we set $F=e^{-i \theta} \tilde{F}$ we get $F^{*} \alpha=$ $* d \psi$. We can make a similar choice for $\Phi$ and note that since $\Psi \in C^{2}$, standard linear theory for elliptic PDEs gives that our $\Phi$ is in $C^{1}$. Thus we get

$$
F^{*} \alpha-\Phi^{*} \alpha=* d f \quad \text { for some function } f
$$

This implies that

$$
\begin{equation*}
\int_{\mathbf{S}^{2}}\left|F^{*} \alpha-\Phi^{*} \alpha\right|^{2}=\|d * d f\|_{W^{-1,2}}^{2}=\left\|F^{*} d \alpha-\Phi^{*} d \alpha\right\|_{W^{-1,2}}^{2} \tag{43}
\end{equation*}
$$

By (22) we have $2 F^{*} d \alpha=M^{*} \pi^{*}\left(i d z_{1} \wedge d \bar{z}_{1}+i d z_{2} \wedge d \bar{z}_{2}\right)=2 M^{*} \omega$ and $2 \Phi^{*} d \alpha=2 \Psi^{*} \omega$. Thus, we conclude that

$$
\begin{align*}
\||d F|-|d \Phi|\|_{L^{2}} & \leq\left\|\left|d_{1} F\right|-\left|d_{1} \Phi\right|\right\|_{L^{2}}+\left\|\left|d_{2} F\right|-\left|d_{2} \Phi\right|\right\|_{L^{2}} \\
& =\left\|M^{*} \omega-\Psi^{*} \omega\right\|_{W^{-1,2}}+\||d M|-|d \Psi|\|_{L^{2}} \tag{44}
\end{align*}
$$

Second Step In this step we show how to estimate $\left\|M^{*} \omega-\Psi^{*} \omega\right\|_{W^{-1,2}}$.
Recall that $\Psi\left(\mathbf{S}^{2}\right) \subset \tilde{D}$, which is the geodesic disk $D_{\pi / 2+1}(p)$. Denote by $n$ the antipodal of $p$. From the area formula there exists a constant $C_{1}$ such that

$$
\operatorname{ar}\left(M\left(\mathbf{S}^{2}\right) \cap D_{1 / 2}(n)\right) \leq C_{1}\|M-\Psi\|_{L^{2}\left(\mathbf{S}^{2}\right)}^{2}
$$

Therefore if $\|M-\Psi\|_{L^{2}\left(\mathbf{S}^{2}\right)}^{2}$ is sufficiently small, $\operatorname{ar}\left(D_{1 / 2}(n) \backslash M\left(\mathbf{S}^{2}\right)\right) \geq C_{2}$ for some positive constant $C_{2}$. We claim the existence of a 1-form $\eta$ such that:

$$
-\omega=d \eta \text { on } \tilde{D} \cup M\left(\mathbf{S}^{2}\right)
$$

$-\|\eta\|_{L^{\infty}} \leq \frac{C}{C_{2}}$;
$-|\eta(x)-\eta(y)| \leq C|x-y|$ for every $x, y \in \tilde{D} ;$
where $C$ is a universal constant.
We construct $\eta$ in the following way. First, for every $x \in \mathbf{S}^{2}$ we take the form $\eta_{x} \in C^{\infty}\left(\mathbf{S}^{2} \backslash\{x\}\right) \cap L^{1}\left(\mathbf{S}^{2}\right)$ defined in 3.5.1 of [4]. This "canonical" form has a singularity in $x$ but satisfies $d \eta_{x}=\omega$ on $\mathbf{S}^{2} \backslash\{x\}$.

Then we take a closed set $E \subset D_{1 / 2}(n) \backslash M\left(\mathbf{S}^{2}\right)$ such that

$$
\operatorname{ar}(E)=\frac{1}{2} \operatorname{ar}\left(D_{1 / 2}(n) \backslash M\left(\mathbf{S}^{2}\right)\right)
$$

and we define

$$
\eta:=\frac{1}{\operatorname{ar}(E)} \int_{x \in E} \eta_{x}
$$

Clearly $d \eta=\omega$ on $D \cup M\left(\mathbf{S}^{2}\right) \subset \mathbf{S}^{2} \backslash E$. Moreover, $\eta$ is smooth on the closure of $\tilde{D}$. The estimate $\|\eta\|_{L^{\infty}} \leq C(\operatorname{ar}(E))^{-1}$ can be proved as in 3.5.5 of [4]. Finally we compute

$$
\begin{aligned}
\left\|M^{*} \omega-\Psi^{*} \omega\right\|_{W^{-1,2}\left(\mathbf{S}^{2}\right)} & =\left\|d\left(M^{*} \eta-\Psi^{*} \eta\right)\right\|_{W^{-1,2}} \\
& =\sup _{\|\varphi\|_{W^{1,2}}=1} \int_{\mathbf{S}^{2}} \varphi d\left(M^{*} \eta-\Psi^{*} \eta\right) \\
& =\sup _{\|\varphi\|_{W^{1,2}}=1} \int_{\mathbf{S}^{2}} d \varphi \wedge\left(M^{*} \eta-\Psi^{*} \eta\right) .
\end{aligned}
$$

Now, write $\eta=f_{1} d x_{1}+f_{2} d x_{2}$ in some local coordinates and note that

$$
\begin{align*}
\Psi^{*}\left(f_{i} d x_{i}\right)-M^{*}\left(f_{i} d x_{i}\right) & =f_{i}(\Psi) d \Psi_{i}-f_{i}(M) d M_{i} \\
& =\left[f_{i}(\Psi)-f_{i}(M)\right] d \Psi_{i}+f_{i}(M) d\left[\Psi_{i}-M_{i}\right] . \tag{45}
\end{align*}
$$

Set $\mathbf{S}^{b}:=\{p \mid M(p) \notin \tilde{D}\}$ and $\mathbf{S}^{g}:=\mathbf{S}^{2} \backslash \mathbf{S}^{b}$. Then we have

$$
\left|\Psi^{*} \eta-M^{*} \eta\right| \leq \begin{cases}C|d \Psi||\Psi-M|+C|d \Psi-d M| & \text { on } \mathbf{S}^{g} \\ 2 C|d \Psi|+C|d \Psi-d M| & \text { on } \mathbf{S}^{b} .\end{cases}
$$

Thus we can estimate

$$
\begin{aligned}
& \left|\int_{\mathbf{S}^{2}} d \varphi \wedge\left(\Psi^{*} \eta-M^{*} \eta\right)\right| \\
& \quad \leq C\|\Psi\|_{C^{1}} \int_{\mathbf{S}^{8}}|d \varphi||\Psi-M|+2 C\|\Psi\|_{C^{1}} \int_{\mathbf{S}^{b}}|d \varphi|+C \int_{\mathbf{S}^{2}}|d \varphi \| d(\Psi-M)| \\
& \leq C\|\Psi\|_{C^{1}}\|d \varphi\|_{L^{2}}\|\Psi-M\|_{L^{2}}+2 C\|\Psi\|_{C^{1}}\|d \varphi\|_{L^{2}}\left(\operatorname{ar}\left(\mathbf{S}^{b}\right)\right)^{1 / 2} \\
& \quad+C\|d \varphi\|_{L^{2}}\|d \Psi-d M\|_{L^{2}} .
\end{aligned}
$$

Recalling that $\|d \varphi\|_{L^{2}} \leq\|\varphi\|_{W^{1,2}}=1$ and that $\left(\operatorname{ar}\left(\mathbf{S}^{b}\right)\right)^{1 / 2} \leq C\|\Psi-M\|_{L^{2}}$, we derive

$$
\begin{equation*}
\left\|M^{*} \omega-\Psi^{*} \omega\right\|_{W^{-1,2}} \leq C_{1}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} . \tag{46}
\end{equation*}
$$

This, together with (44), gives

$$
\begin{equation*}
\||d F|-|d \Phi|\|_{L^{2}} \leq C_{2}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} . \tag{47}
\end{equation*}
$$

Moreover, for a later use, we remark that (46) and (43) give

$$
\begin{equation*}
\left\|F^{*} \alpha-\Phi^{*} \alpha\right\|_{L^{2}\left(\mathbf{S}^{2}\right)} \leq C_{3}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)} . \tag{48}
\end{equation*}
$$

Third Step We now come to the proof of (26).
We begin with the following pointwise inequality:
$|d F-d \Phi|_{\mathcal{L}}^{2} \leq C_{2}(|d F|+|d \Phi|)^{2}|F-\Phi|^{2}+4|d M-d \Phi|^{2}+4\left|F^{*} \alpha-\Phi^{*} \alpha\right|^{2}$,
where $\alpha$ is the differential form $-i \bar{z}_{1} d z_{2}-i \bar{z}_{2} d z_{2}$, which satisfies (42).
In order to prove (49), for every $\xi \in T_{a} \mathbf{S}^{3}, \zeta \in T_{b} \mathbf{S}^{3}$ we define a distance $d(\xi, \zeta)$ in the following way. We write $\xi=\tilde{\xi}+t i a$ and $\zeta=\tilde{\zeta}+\tau i b$, where $\tilde{\xi} \in T N_{a}, \tilde{\zeta} \in T N_{b}$ and $t, \tau \in \mathbf{R}$ (see (41)). Then we set

$$
d(\xi, \eta):=\sqrt{\left|d \pi_{a}(\xi)-d \pi_{b}(\zeta)\right|^{2}+|\tau-t|^{2}} .
$$

Now, construct the function $f: T_{a} \mathbf{S}^{3} \times T_{b} \mathbf{S}^{3} \rightarrow \mathbf{R}$ given by

$$
f(a, b, \xi, \eta)=\| \xi-\left.\eta\right|_{\mathcal{L}}-d(\xi, \eta) \mid .
$$

Note that both $d$ and $|\cdot|_{\mathcal{L}}$ are locally Lipschitz in $a, b, \xi$, and $\zeta$. Moreover

$$
d(\xi, \eta)=|\xi-\eta|=|\xi-\eta|_{\mathcal{L}} \quad \text { for } \xi, \eta \in T_{a} \mathbf{S}^{3},
$$

which translates into $f(a, a, \xi, \eta)=0$. This condition and the locally Lipschitz property of $f$ gives the existence of a constant $C$ such that:

$$
\begin{equation*}
f(a, b, \xi, \eta) \leq C|a-b| \quad \text { for }|\xi|+|\eta| \leq 2 \tag{50}
\end{equation*}
$$

Given any $\xi, \eta$ we define $M:=\max \{|\xi|,|\zeta|\}$ and $\hat{\xi}:=\xi / M, \hat{\zeta}:=\zeta / M$. Then we can compute

$$
\begin{align*}
|\xi-\eta|_{\mathcal{L}} & =M|\hat{\xi}-\hat{\zeta}|_{\mathcal{L}} \leq M d(\hat{\xi}, \hat{\zeta})+C M|a-b| \\
& \leq d(\xi, \eta)+C(|\xi|+|\eta|)|a-b| \tag{51}
\end{align*}
$$

From this we easily get (49). Integrating (49) and recalling (48) we get the inequality

$$
\begin{align*}
& \|d F-d \Phi\| \|_{L^{2}\left(D_{r}\right)}^{2} \\
& \quad \leq C_{1} \int_{D_{r}}(|d F|+|d \Phi|)^{2}|F-\Phi|^{2}+C_{2}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}^{2} \tag{52}
\end{align*}
$$

Moreover, we have

$$
\begin{align*}
& \int_{D_{r}}(|d F|+|d \Phi|)^{2}|F-\Phi|^{2} \leq \int_{D_{r}}\left(8|d \Phi|^{2}+2| | d F\left|-|d \Phi|^{2}\right)|F-\Phi|^{2}\right. \\
& \quad \leq 4\||d F|-|d \Phi|\|_{L^{2}\left(D_{r}\right)}^{2}+8\|\Phi\|_{C^{1}}^{2}\|F-\Phi\|_{L^{2}\left(D_{r}\right)}^{2} \\
& \quad \stackrel{(47)}{\leq} C_{3}\|M-\Psi\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}^{2}+C_{4}\|F-\Phi\|_{L^{2}\left(D_{r}\right)}^{2} \tag{53}
\end{align*}
$$

Plugging (53) into (52) we derive

$$
\begin{equation*}
\|d F-d \Phi\|\left\|_{L^{2}\left(D_{r}\right)} \leq C_{5}\right\| M-\Psi\left\|_{W^{1,2}\left(\mathbf{S}^{2}\right)}+C_{6}\right\| F-\Phi \|_{L^{2}\left(D_{r}\right)} \tag{54}
\end{equation*}
$$

Given $\theta \in \mathbf{R}$, define $\tilde{\Phi}=e^{i \theta} \Phi$. Then, clearly $|d \Phi-d \tilde{\Phi}|_{\mathcal{L}}=0$. Note that $\tilde{\Phi}$ is a lifting of $\Psi$ and that all the estimates derived for $\Phi$ holds for $\tilde{\Phi}$ as well. Hence from (54) we get (26).

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