

Lecture No 3

Fast Diffusion

The Ricci flow on surfaces

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We will discuss non-linear parabolic equations of **fast diffusion**.
Our model is the **fast diffusion** equation

$$u_t = \Delta u^m = \operatorname{div}(m u^{m-1} \nabla u), \quad m < 1.$$

- It appears in physical applications such as diffusion in plasma, thin liquid film dynamics.
- The case $m = 0$ in dimension $n = 2$ corresponds to the **Ricci flow** on surfaces.
- The case $m = \frac{n-2}{n+2}$ in dimensions $n \geq 3$ corresponds to the **Yamabe flow**.
- Since, the diffusivity $D(u) = m u^{m-1} \uparrow +\infty$, as $u \downarrow 0$ the equation becomes **singular** at $u = 0$, resulting to the phenomenon of **fast-diffusion**.

Scaling and the Barenblatt solution

Scaling: If u solves the fast diffusion equation $u_t = \Delta u^m$, then

$$\tilde{u}(x, t) = \gamma^{-1} u(\alpha x, \beta t), \quad \gamma = \left(\frac{\beta}{\alpha^2} \right)^{\frac{1}{1-m}}$$

also solves the same equation.

Self-Similar solution: There exists a self-similar solution

$$U(x, t) = t^{-\lambda} \left(C + k \frac{|x|^2}{t^{2\mu}} \right)^{-\frac{1}{1-m}}$$

with

$$\lambda^{-1} = \frac{2}{n} - (1 - m), \quad \mu = \frac{\lambda}{n}, \quad k = \frac{\lambda(1 - m)}{2mn}.$$

The above is a solution if $\frac{2}{n} - (1 - m) > 0$, i.e. $m > \frac{n-2}{n}$.

The exponent $m = \frac{n-2}{n}$ is **critical**.

The Aronson-Bénilan inequality

Every solution u to the fast-diffusion equation

$$u_t = \Delta u^m, \quad \frac{(n-2)_+}{n} < m < 1$$

satisfies the differential inequality

$$(*) \quad u_t \geq -\frac{\lambda u}{t}$$

with $\lambda^{-1} = \frac{2}{n} - (1-m) > 0$ iff $m > \frac{(n-2)_+}{n}$.

The *pressure* $v := \frac{m}{1-m} u^{-(1-m)}$ which evolves by the equation

$$v_t = (1-m)v\Delta v - |\nabla v|^2$$

satisfies the **sharp** differential inequality

$$(**) \quad \Delta v \leq \frac{\lambda}{t}.$$

Remark: The Aronson-Bénilan (*) inequality follows from (**).
The differential inequality (**) becomes an **equality** when v is the Barenblatt solution.

The Li-Yau type Harnack Inequality

Combining the Aronson-Bénilan inequality $\Delta v \leq \frac{\lambda}{t}$ with the evolution equation for v gives the following **Li-Yau type** differential inequality:

$$-v_t + (1 - m) \lambda \frac{v}{t} \geq |\nabla v|^2.$$

Integrating this inequality on optimal paths we obtain:

Harnack Inequality (Auchmuty-Bao and Hamilton)

If $0 < t_1 < t_2$, then

$$v(x_2, t_2) \leq \left(\frac{t_2}{t_1}\right)^\mu \left[v(x_1, t_1) + \frac{\delta}{4} \frac{|x_2 - x_1|^2}{t_2^\delta - t_1^\delta} t_1^\mu \right].$$

with $\mu = (1 - m) \lambda > 0$ and $\delta = \frac{2\lambda}{n}$.

Application: Solutions of $u_t = \Delta u^m$ with $\frac{(n-2)_+}{n} < m < 1$ satisfy the lower bound:

$$u(x, t) \geq c(t) (1 + |x|^2)^{-\frac{1}{1-m}}.$$

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The other Aronson-Bénilan inequality

A simple scaling argument shows that every solution u to the fast-diffusion equation $u_t = \Delta u^m$, $0 \leq m < 1$ satisfies the differential inequality

$$(*) \quad u_t \leq \frac{u}{(1-m)t}.$$

Integrating $(*)$ on a time interval $[t_1, t_2]$ gives:

$$u(x, t_2) \leq u(x, t_1) \left(\frac{t_2}{t_1} \right)^{\frac{1}{1-m}}, \quad \forall x$$

i.e. the L^∞ norm of a solution doesn't blow up.

Remark: In the range of exponents $\frac{(n-2)_+}{n} < m < 1$, solutions u exhibit a **regularizing effect** from L^1_{loc} to L^∞_{loc} :

$$\sup_{|x| \leq R} u(x, t) \leq F \left(t, R, \int_{B_{2R}} u_0(x) dx \right).$$

The Cauchy problem in the super-critical case

Consider the **fast-diffusion** equation

$$(*) \quad u_t = \Delta u^m, \quad \frac{(n-2)_+}{n} < m < 1.$$

In the super-critical case **no growth conditions** need to be imposed on the initial data for existence. More precisely:

- For any nonnegative continuous weak solution u of $(*)$, there exists a unique locally finite Borel measure μ_0 on \mathbb{R}^n such that

$$\lim_{t \downarrow 0} u(\cdot, t) = \mu_0 \quad \text{in } D'(\mathbb{R}^n).$$

- The trace μ_0 determines the solution **uniquely**.
- For any locally finite Borel measure μ_0 on \mathbb{R}^n there exists a continuous weak solution u of $(*)$ in $S_\infty = \mathbb{R}^n \times (0, \infty)$ with initial trace μ_0 .
- **Remark:** Solutions exist **for all time** $t > 0$ independently of their initial data.

The sub-critical case $m < (n - 2)_+/n$

- In the sub-critical case $m < \frac{(n-2)_+}{n}$ the analogues of the above results **do not hold true**. In particular, there exists **no solution** with initial data the **Dirac mass**.
- One particularly interesting case of **critical** fast diffusion is the case where $u_t = \Delta \log u$, corresponding to the limiting case of $u_t = \frac{\Delta u^m}{m} = \operatorname{div}(u^{-1} \nabla u)$, when $m \rightarrow 0$. In **dim $n = 2$** it corresponds to the **Ricci flow** on surfaces.
- The sub-critical fast-diffusion equation $u_t = \Delta u^{\frac{n-2}{n+2}}$ corresponds to the evolution of a conformally flat metric $g_{ij} = u^{\frac{4}{n+2}} dx_i dx_j$ by the **Yamabe flow** $(g_{ij})_t = -R g_{ij}$, where $R = -u^{-1} \Delta u^{\frac{n-2}{n+2}}$ denotes the **scalar curvature** of g .
- The Yamabe flow was used to give an alternative proof of the well known **Yamabe conjecture**.

Logarithmic fast-diffusion

We will now consider the **logarithmic fast-diffusion** equation

$$(*) \quad u_t = \Delta \log u, \quad \text{in } \mathbb{R}^n \times [0, T), \quad T > 0.$$

- Lions and Toscani: $(*)$ arises as a singular limit for finite velocity **Boltzmann kinetic models**.
- Kurtz: $(*)$ describes the **limiting density distribution** of two gases moving against each other and obeying the Boltzmann equation.
- In dimension $n = 2$ equation $(*)$ arises as a model for **long Va-der-Wals** interactions in **thin films** of a fluid spreading on a solid surface, if certain nonlinear fourth order effects are neglected.
- Angenent and L. Wu: In dim $n = 2$, equation $(*)$ represents the evolution of the conformal metric $g_{ij} = u dx_i dx_j$ under the **Ricci flow**, $(g_{ij})_t = -2 R_{ij}$.

The Ricci flow on \mathbb{R}^2

- In 1982 **R. Hamilton** introduced the **Ricci flow**, namely the evolution of a Riemannian metric g_{ij} by

$$(RF) \quad \frac{\partial g_{ij}}{\partial t} = -2 R_{ij}$$

where R_{ij} denotes the **Ricci curvature** of the metric g_{ij} .

- In **dimension $n = 2$** we have:

$$g_{ij} = u dx_i dx_j, \quad R_{ij} = \frac{1}{2} R g_{ij}, \quad R = -\frac{\Delta \log u}{u}$$

where R denotes the scalar curvature of g_{ij} .

- Hence, in **dimension $n = 2$** the Ricci flow (RF) is equivalent to the equation:

$$(*) \quad u_t = \Delta \log u.$$

Examples of solutions of $u_t = \Delta \log u$ on \mathbb{R}^2

- **Contracting spheres:** $u(x, t) = \frac{8\lambda(T-t)}{(\lambda+|x|^2)^2}$, $\lambda > 0$.

They are *ancient solutions* which vanish at time $t = T$ and:

$$\frac{d}{dt} \int_{\mathbb{R}^2} u \, dx = \int_{\mathbb{R}^2} R u = -4\pi.$$

- **Cigar solution:** $u(x, t) = \frac{1}{\lambda|x|^2 + e^{4\lambda t}}$.

They are *eternal complete non-compact* solutions which look like **cigars** and have infinite area, i.e.

$$\int_{\mathbb{R}^2} u \, dx = \infty.$$

- **Cusp solution:** $u(x, t) = \frac{2t}{|x|^2 \log^2 |x|}$, $|x| > 1$.

They are *complete non-compact* solutions which look like **cusps** and have finite area.

Solutions on orbifolds

- **Cylindrical change of coordinates:**

If $u(r, t)$ is a **radial** solution of $u_t = \Delta \log u$ in $\mathbb{R}^2 \times (0, T)$, then $v(s, t) := r^2 u(r, t)$, $s = \log r$ is a solution of

$$(*) \quad v_t = (\log v)_{ss}, \quad \text{on } \mathbb{R} \times (0, T).$$

- We seek for solutions $v^\mu(s, t)$ of $(*)$ with $T = 1$ in the form $v^\mu(s, t) = (1 - t) g(s - \gamma_\mu \log(1 - t))$. It follows that for each $\mu > 0$, there exists γ_μ and solution v^μ such that

$$\frac{d}{dt} \int_{-\infty}^{\infty} v^\mu(s, t) ds = -(2 + \mu).$$

Then, $u^\mu := |x|^{-2} v(\log |x|, t)$ defines a radial solution of $u_t = \Delta \log u$ on in $\mathbb{R}^2 \times (0, T)$ such that

$$\frac{d}{dt} \int_{\mathbb{R}^2} u^\mu(x, t) dx = -2\pi (2 + \mu).$$

The u^μ , $\mu \neq 2$ define metrics on **orbifolds**. When $\mu = 2$ we recover the **contracting spheres**.

The Cauchy problem

Consider the **Cauchy problem**

$$(*) \quad \begin{cases} u_t = \Delta \log u & \text{in } \mathbb{R}^2 \times [0, T) \\ u(x, 0) = f & \text{on } \mathbb{R}^2 \end{cases}$$

with initial data $f \geq 0$.

- If $\int_{\mathbb{R}^2} f \, dx < \infty$, then $\forall \mu \geq 0, \exists u_\mu$ solution of $(*)$ on $\mathbb{R}^2 \times (0, T_\mu)$ with $T_\mu = \frac{1}{2\pi(2+\mu)} \int_{\mathbb{R}^2} f(x) \, dx$ satisfying

$$\frac{d}{dt} \int_{\mathbb{R}^2} u^\mu(x, t) \, dx = -2\pi(2 + \mu).$$

- If $\int_{\mathbb{R}^2} f \, dx = \infty$, $\exists u$ solution of $(*)$ on $\mathbb{R}^2 \times (0, \infty)$.

Remarks on the Cauchy problem

- If $\int_{\mathbb{R}^2} f \, dx < \infty$, then every solution vanishes at time $T \leq T_{\max}$, with $T_{\max} = \frac{1}{4\pi} \int_{\mathbb{R}^2} f(x) \, dx$.
- The maximal solution ($\mu = 0$) defines **complete** metrics on \mathbb{R}^2 of **finite area** which evolve by the Ricci flow and behave as **cusps** near infinity.
- The intermediate solution u_μ with $\mu = 2$ corresponds to **smooth** metrics on S^2 evolving by the Ricci flow.
- All other solutions u_μ ($\mu \neq 0, 2$) correspond to metrics on **orbifolds** evolving by the Ricci flow.
- Radial u_μ , $\mu > 0$ are characterized by the **outgoing flux** at infinity:

$$\lim_{r \rightarrow \infty} r (\log u_\mu)_r = -(2 + \mu).$$

- **Conclusion:** A strong **non-uniqueness** phenomenon takes place. This is in contrast to the **uniqueness** of solutions to equations $u_t = \Delta u^m$, $m > 0$.

The Ricci flow on surfaces

- Consider the intermediate solution u of $u_t = \Delta \log u$ with area decay: $\frac{d}{dt} \int_{\mathbb{R}^2} u(x, t) dx = -8\pi$.
- An example of such a solution is the **contracting spheres**

$$u_s(x, t) = \frac{8(T - t)}{(1 + |x|^2)^2}.$$

- Any such solution can be lifted on S^2 defining **smooth** metrics $ds^2 = \bar{u} ds_p^2$ evolving by the Ricci flow. This solution will exist up to time $T = \frac{1}{8\pi} \int_{S^2} \bar{u}(\cdot, 0) da$.
- The special solution u_s when lifted on S^2 is $\bar{u}_s = 2(T - t)$.
- **Hamilton:** Any solution $ds^2 = \bar{u} ds_p^2$ of the Ricci flow on S^2 will converge (after re-scaling) as $t \rightarrow T$ to the round sphere solution \bar{u}_s .

Vanishing behavior of the maximal solution

- Consider the **maximal solution** u of $u_t = \Delta \log u$ with compactly supported initial data u_0 . Its area decays as:

$$\frac{d}{dt} \int_{\mathbb{R}^2} u(x, t) dx = -4\pi.$$

- Hence u will **vanish** at time $T = \frac{1}{4\pi} \int_{\mathbb{R}^2} u_0(x) dx$.
- Any such solution will enjoy the asymptotic behavior:

$$u(x, t) = \frac{2t}{|x|^2 \log^2 |x|} (1 + o(1)), \quad \forall t < T$$

however the bound deteriorates as $t \rightarrow T$.

- Question:** What is the **vanishing behavior** of u as $t \rightarrow T$?

Vanishing behavior of the maximal solution

Jointly with M. del Pino and N. Sesum we have shown the vanishing asymptotic behavior for the maximal solution u :

- On the **outer region**: $(T - t) \log |x| > T$, we have

$$u(x, t) \approx \frac{2T}{|x|^2 \log^2 |x|}, \quad \text{as } t \rightarrow T^-.$$

- On the **inner region**: $(T - t) \log |x| > T$, or $|x| = O(e^{\frac{T}{T-t}})$, u has the **self-similar** profile:

$$u(x, t) \approx (T - t)^2 e^{-\frac{2T}{T-t}} \phi(e^{-\frac{T}{T-t}} |x|)$$

where $\phi(r) = \frac{2T^{-1}}{r^2 + b}$.

This shows that in the inner region, after re-scaling, the solution u converges to one of the **cigar** solutions.

- **Remark**: The vanishing of the maximum solutions if of **type II**.

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